

# Solving Nonlinear Wave Equations and Lattices with Mathematica

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## OUTLINE

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## Purpose & Motivation

- Develop various symbolic algorithms to compute exact solutions of nonlinear (systems) of partial differential equations (PDEs) and differential-difference equations (DDEs, lattices).
- Fully **automate** the tanh and sech methods to compute closed form solitary wave solutions.
- Solutions of tanh or sech type **model** solitary waves in fluid dynamics, plasmas, electrical circuits, optical fibers, bio-genetics, etc.
- Class of nonlinear PDEs and DDEs solvable with the tanh/sech method includes famous evolution and wave equations.  
Typical examples: Korteweg-de Vries, Fisher and Boussinesq PDEs, Toda and Volterra lattices (DDEs).
- **Research aspect**: Design a high-quality application package for the computation of exact solitary wave solutions of large classes of nonlinear evolution and wave equations.
- **Educational aspect**: Software as course ware for courses in nonlinear PDEs, theory of nonlinear waves, integrability, dynamical systems, and modeling with symbolic software.
- Users: scientists working on nonlinear wave phenomena in fluid dynamics, nonlinear networks, elastic media, chemical kinetics, material science, bio-sciences, plasma physics, and nonlinear optics.

## Typical Examples of Single PDEs and Systems of PDEs

- The Korteweg-de Vries (KdV) equation:

$$u_t + \alpha u u_x + u_{3x} = 0.$$

Solitary wave solution:

$$u(x, t) = \frac{8c_1^3 - c_2}{6\alpha c_1} - \frac{2c_1^2}{\alpha} \tanh^2 [c_1 x + c_2 t + \delta],$$

or, equivalently,

$$u(x, t) = -\frac{4c_1^3 + c_2}{6\alpha c_1} + \frac{2c_1^2}{\alpha} \operatorname{sech}^2 [c_1 x + c_2 t + \delta].$$

- The modified Korteweg-de Vries (mKdV) equation:

$$u_t + \alpha u^2 u_x + u_{3x} = 0.$$

Solitary wave solution:

$$u(x, t) = \pm \sqrt{\frac{6}{\alpha}} c_1 \operatorname{sech} [c_1 x - c_1^3 t + \delta].$$

- The Fisher equation:

$$u_t - u_{xx} - u(1 - u) = 0.$$

Solitary wave solution:

$$u(x, t) = \frac{1}{4} \pm \frac{1}{2} \tanh \xi + \frac{1}{4} \tanh^2 \xi,$$

with

$$\xi = \pm \frac{1}{2\sqrt{6}} x \pm \frac{5}{12} t + \delta.$$

- The generalized Kuramoto-Sivashinski equation:

$$u_t + uu_x + u_{xx} + \sigma u_{3x} + u_{4x} = 0.$$

Solitary wave solutions

(ignoring symmetry  $u \rightarrow -u, x \rightarrow -x, \sigma \rightarrow -\sigma$ ):

For  $\sigma = 4$ :

$$u(x, t) = 9 - 2c_2 - 15 \tanh\xi (1 + \tanh\xi - \tanh^2\xi),$$

with  $\xi = \frac{x}{2} + c_2 t + \delta$ .

For  $\sigma = 0$ :

$$u(x, t) = -2\sqrt{\frac{19}{11}}c_2 - \frac{135}{19}\sqrt{\frac{11}{19}}\tanh\xi + \frac{165}{19}\sqrt{\frac{11}{19}}\tanh^3\xi,$$

with  $\xi = \frac{1}{2}\sqrt{\frac{11}{19}}x + c_2 t + \delta$ .

For  $\sigma = \frac{12}{\sqrt{47}}$ :

$$\begin{aligned} u(x, t) &= \frac{45 \mp 4418c_2}{47\sqrt{47}} \pm \frac{45}{47\sqrt{47}} \tanh\xi \\ &\quad - \frac{45}{47\sqrt{47}} \tanh^2\xi \pm \frac{15}{47\sqrt{47}} \tanh^3\xi, \end{aligned}$$

with  $\xi = \pm \frac{1}{2\sqrt{47}}x + c_2 t + \delta$ .

For  $\sigma = 16/\sqrt{73}$ :

$$\begin{aligned} u(x, t) &= \frac{2(30 \mp 5329c_2)}{73\sqrt{73}} \pm \frac{75}{73\sqrt{73}} \tanh\xi \\ &\quad - \frac{60}{73\sqrt{73}} \tanh^2\xi \pm \frac{15}{73\sqrt{73}} \tanh^3\xi, \end{aligned}$$

with  $\xi = \pm \frac{1}{2\sqrt{73}}x + c_2 t + \delta$ .

- Three-dimensional modified Korteweg-de Vries equation:

$$u_t + 6u^2u_x + u_{xyz} = 0.$$

Solitary wave solution:

$$u(x, y, z, t) = \pm \sqrt{c_2 c_3} \operatorname{sech} [c_1 x + c_2 y + c_3 z - c_1 c_2 c_3 t + \delta].$$

- The Boussinesq (wave) equation:

$$u_{tt} - \beta u_{2x} + 3uu_{2x} + 3u_x^2 + \alpha u_{4x} = 0,$$

**or** written as a first-order system ( $v$  auxiliary variable):

$$\begin{aligned} u_t + v_x &= 0, \\ v_t + \beta u_x - 3uv_x - \alpha u_{3x} &= 0. \end{aligned}$$

Solitary wave solution:

$$\begin{aligned} u(x, t) &= \frac{\beta c_1^2 - c_2^2 + 8\alpha c_1^4}{3c_1^2} - 4\alpha c_1^2 \tanh^2 [c_1 x + c_2 t + \delta] \\ v(x, t) &= b_0 + 4\alpha c_1 c_2 \tanh^2 [c_1 x + c_2 t + \delta]. \end{aligned}$$

- The Broer-Kaup system:

$$\begin{aligned} u_{ty} + 2(uu_x)_y + 2v_{xx} - u_{xxy} &= 0, \\ v_t + 2(uv)_x + v_{xx} &= 0. \end{aligned}$$

Solitary wave solution:

$$\begin{aligned} u(x, t) &= -\frac{c_3}{2c_1} + c_1 \tanh [c_1 x + c_2 y + c_3 t + \delta] \\ v(x, t) &= c_1 c_2 - c_1 c_2 \tanh^2 [c_1 x + c_2 y + c_3 t + \delta] \end{aligned}$$

## Typical Examples of DDEs (lattices)

- The Toda lattice:

$$\ddot{u}_n = (1 + \dot{u}_n) (u_{n-1} - 2u_n + u_{n+1}).$$

Solitary wave solution:

$$u_n(t) = a_0 \pm \sinh(c_1) \tanh [c_1 n \pm \sinh(c_1) t + \delta].$$

- The Volterra lattice:

$$\begin{aligned}\dot{u}_n &= u_n(v_n - v_{n-1}) \\ \dot{v}_n &= v_n(u_{n+1} - u_n).\end{aligned}$$

Solitary wave solution:

$$\begin{aligned}u_n(t) &= -c_2 \coth(c_1) + c_2 \tanh [c_1 n + c_2 t + \delta] \\ v_n(t) &= -c_2 \coth(c_1) - c_2 \tanh [c_1 n + c_2 t + \delta].\end{aligned}$$

- The Relativistic Toda lattice:

$$\begin{aligned}\dot{u}_n &= (1 + \alpha u_n)(v_n - v_{n-1}) \\ \dot{v}_n &= v_n(u_{n+1} - u_n + \alpha v_{n+1} - \alpha v_{n-1}).\end{aligned}$$

Solitary wave solution:

$$\begin{aligned}u_n(t) &= -c_2 \coth(c_1) - \frac{1}{\alpha} + c_2 \tanh [c_1 n + c_2 t + \delta] \\ v_n(t) &= \frac{c_2 \coth(c_1)}{\alpha} - \frac{c_2}{\alpha} \tanh [c_1 n + c_2 t + \delta].\end{aligned}$$

## Algorithm for Tanh Solutions for PDE system

Given is a system of PDEs of order  $n$

$$\Delta(\mathbf{u}(\mathbf{x}), \mathbf{u}'(\mathbf{x}), \mathbf{u}''(\mathbf{x}), \dots, \mathbf{u}^{(n)}(\mathbf{x})) = \mathbf{0}.$$

Dependent variable  $\mathbf{u}$  has components  $u_i$  (or  $u, v, w, \dots$ )

Independent variable  $\mathbf{x}$  has components  $x_i$  (or  $x, y, z, t$ )

### Step 1:

- Seek solution of the form  $u_i(\mathbf{x}) = U_i(T)$ , with

$$T = \tanh [c_1x + c_2y + c_3z + c_4t] = \tanh \xi.$$

- Observe  $\cosh^2 \xi - \sinh^2 \xi = 1$ ,  $(\tanh \xi)' = 1 - \tanh^2 \xi$   
or  $T' = 1 - T^2$ .
- Repeatedly apply the operator rule

$$\frac{\partial \bullet}{\partial x_i} \rightarrow c_i(1 - T^2) \frac{d \bullet}{dT}$$

This produces a coupled system of Legendre equations of type

$$\mathbf{P}(T, U_i, U_i', \dots, U_i^{(n)}) = 0$$

for  $U_i(T)$ .

- Example: For Boussinesq system

$$u_t + v_x = 0$$

$$v_t + \beta u_x - 3uu_x - \alpha u_{3x} = 0,$$

we obtain after cancelling common factors  $1 - T^2$

$$c_2U' + c_1V' = 0$$

$$\begin{aligned} &c_2V' + \beta c_1U' - 3c_1UU' \\ &+ \alpha c_1^3 [2(1 - 3T^2)U' + 6T(1 - T^2)U'' - (1 - T^2)^2U'''] = 0 \end{aligned}$$

## Step 2:

- Seek polynomial solutions

$$U_i(T) = \sum_{j=0}^{M_i} a_{ij} T^j$$

Balance the highest power terms in  $T$  to determine  $M_i$ .

- Example: Powers for Boussinesq system

$$M_1 - 1 = M_2 - 1, \quad 2M_1 - 1 = M_1 + 1$$

gives  $M_1 = M_2 = 2$ .

Hence,  $U_1(T) = a_{10} + a_{11}T + a_{12}T^2$ ,  $U_2(T) = a_{20} + a_{21}T + a_{22}T^2$ .

## Step 3:

- Determine the algebraic system for the unknown coefficients  $a_{ij}$  by balancing the coefficients of the various powers of  $T$ .
- Example: Boussinesq system

$$a_{11} c_1 (3a_{12} + 2\alpha c_1^2) = 0$$

$$a_{12} c_1 (a_{12} + 4\alpha c_1^2) = 0$$

$$a_{21} c_1 + a_{11} c_2 = 0$$

$$a_{22} c_1 + a_{12} c_2 = 0$$

$$\beta a_{11} c_1 - 3a_{10} a_{11} c_1 + 2\alpha a_{11} c_1^3 + a_{21} c_2 = 0$$

$$-3a_{11}^2 c_1 + 2\beta a_{12} c_1 - 6a_{10} a_{12} c_1 + 16\alpha a_{12} c_1^3 + 2a_{22} c_2 = 0.$$

#### Step 4:

- Solve the nonlinear algebraic system with parameters.  
Reject complex solutions? Test the solutions.
- Example: Solution for Boussinesq case

$$a_{10} = \frac{\beta c_1^2 - c_2^2 + 8\alpha c_1^4}{3c_1^2}$$

$$a_{11} = 0$$

$$a_{12} = -4\alpha c_1^2$$

$$a_{20} = \text{free}$$

$$a_{21} = 0$$

$$a_{22} = 4\alpha c_1 c_2.$$

#### Step 5:

- Return to the original variables.  
Test the final solution in the original equations
- Example: Solitary wave solution for Boussinesq system:

$$u(x, t) = \frac{\beta c_1^2 - c_2^2 + 8\alpha c_1^4}{3c_1^2} - 4\alpha c_1^2 \tanh^2 [c_1 x + c_2 t + \delta]$$

$$v(x, t) = a_{20} + 4\alpha c_1 c_2 \tanh^2 [c_1 x + c_2 t + \delta].$$

## Algorithm for Sech Solutions for PDE system

Given is a system of PDEs of order  $n$

$$\Delta(\mathbf{u}(\mathbf{x}), \mathbf{u}'(\mathbf{x}), \mathbf{u}''(\mathbf{x}), \dots, \mathbf{u}^{(n)}(\mathbf{x})) = \mathbf{0}.$$

Dependent variable  $\mathbf{u}$  has components  $u_i$  (or  $u, v, w, \dots$ )

Independent variable  $\mathbf{x}$  has components  $x_i$  (or  $x, y, z, t$ )

### Step 1:

- Seek solution of the form  $u_i(\mathbf{x}) = U_i(S)$ , with

$$S = \operatorname{sech} [c_1x + c_2y + c_3z + c_4t] = \tanh \xi.$$

- Observe  $(\operatorname{sech} \xi)' = -\tanh \xi \operatorname{sech} \xi$  or  $S' = -TS = -\sqrt{1-S^2} S$ .

Also,  $\cosh^2 \xi - \sinh^2 \xi = 1$ , hence,  $T^2 + S^2 = 1$  and  $\frac{dT}{dS} = -\frac{S}{T}$ .

- Repeatedly apply the operator rule

$$\frac{\partial \bullet}{\partial x_i} \rightarrow -c_i \sqrt{1-S^2} S \frac{d\bullet}{dS}$$

This produces a coupled system of Legendre type equations of type

$$\mathbf{P}(S, U_i, U_i', \dots, U_i^{(m)}) + \sqrt{1-S^2} \mathbf{Q}(S, U_i, U_i', \dots, U_i^{(n)}) = 0$$

for  $U_i(S)$ .

For every equation one **must** have  $P_i = 0$  or  $Q_i = 0$ . Only odd derivatives produce the extra factor  $\sqrt{1-S^2}$ .

Conclusion: The total number of derivatives in each term in the given system should be either even or odd. No mismatch is allowed.

- Example: For the 3D mKdV equation

$$u_t + 6u^2u_x + u_{xyz} = 0.$$

we obtain after cancelling a common factor  $-\sqrt{1-S^2}S$

$$c_4U' + 6c_1U^2U' + c_1c_2c_3[(1-6S^2)U' + 3S(1-2S^2)U'' + S^2(1-S^2)U'''] = 0$$

### Step 2:

- Seek polynomial solutions

$$U_i(S) = \sum_{j=0}^{M_i} a_{ij}S^j$$

Balance the highest power terms in  $S$  to determine  $M_i$ .

- Example: Powers for the 3D mKdV case

$$3M_1 - 1 = M_1 + 1$$

gives  $M_1 = 1$ . Hence,  $U(S) = a_{10} + a_{11}S$ .

### Step 3:

- Determine the algebraic system for the unknown coefficients  $a_{ij}$  by balancing the coefficients of the various powers of  $S$ .
- Example: System for 3D mKdV case

$$a_{11}c_1(a_{11}^2 - c_2c_3) = 0$$

$$a_{11}(6a_{10}^2c_1 + c_1c_2c_3 + c_4) = 0$$

$$a_{10}a_{11}^2c_1 = 0$$

#### Step 4:

- Solve the nonlinear algebraic system with parameters.  
Reject complex solutions? Test the solutions.
- Example: Solution for 3D mKdV case

$$\begin{aligned}a_{10} &= 0 \\a_{11} &= \pm\sqrt{c_1 c_3} \\c_4 &= -c_1 c_2 c_3\end{aligned}$$

#### Step 5:

- Return to the original variables.  
Test the final solution in the original equations
- Example: Solitary wave solution for the 3D mKdV equation

$$u(x, y, z, t) = \pm\sqrt{c_2 c_3} \operatorname{sech}(c_1 x + c_2 y + c_3 z - c_1 c_2 c_3 t).$$

## Extension: Tanh Solutions for DDE system

Given is a system of differential-difference equations (DDEs) of order  $n$

$$\Delta(\dots, \mathbf{u}_{n-1}, \mathbf{u}_n, \mathbf{u}_{n+1}, \dots, \dot{\mathbf{u}}_n, \dots, \mathbf{u}_n^{(m)}, \dots) = 0.$$

Dependent variable  $\mathbf{u}_n$  has components  $u_{i,n}$  (or  $u_n, v_n, w_n, \dots$ )

Independent variable  $\mathbf{x}$  has components  $x_i$  (or  $n, t$ ).

No derivatives on shifted variables are allowed!

### Step 1:

- Seek solution of the form  $u_{i,n}(\mathbf{x}) = U_{i,n}(T(n))$ , with

$$T(n) = \tanh [c_1 n + c_2 t + \delta] = \tanh \xi.$$

- Note that the argument  $T$  depends on  $n$ . Complicates matters.
- Repeatedly apply the operator rule on  $u_{i,n}$

$$\frac{\partial \bullet}{\partial t} \rightarrow c_2(1 - T^2) \frac{d \bullet}{dT}$$

This produces a coupled system of Legendre equations of type

$$\mathbf{P}(T, U_{i,n}, U'_{i,n}, \dots) = 0$$

for  $U_{i,n}(T)$ .

- Example: Toda lattice

$$\ddot{u}_n = (1 + \dot{u}_n)(u_{n-1} - 2u_n + u_{n+1}).$$

transforms into

$$\begin{aligned} & c_2^2(1 - T^2) [2TU'_n - (1 - T^2)U''_n] \\ & + [1 + c_2(1 - T^2)U'_n] [U_{n-1} - 2U_n + U_{n+1}] = 0 \end{aligned}$$

## Step 2:

- Seek polynomial solutions

$$U_{i,n}(T(n)) = \sum_{j=0}^{M_i} a_{ij} T(n)^j$$

For  $U_{n+p}$ ,  $p \neq 0$ , there is a phase shift:

$$U_{i,n\pm p}(T(n \pm p)) = \sum_{j=0}^{M_i} a_{i,j} [T(n \pm p)]^j = \sum_{j=0}^{M_i} a_{i,j} \left[ \frac{T(n) \pm \tanh(pc_1)}{1 \pm T(n) \tanh(pc_1)} \right]^j$$

Balance the highest power terms in  $T(n)$  to determine  $M_i$ .

- Example: Powers for Toda lattice

$$2M_1 - 1 = M_1 + 1$$

gives  $M_1 = 1$ .

Hence,

$$\begin{aligned} U_n(T(n)) &= a_{10} + a_{11}T(n) \\ U_{n-1}(T(n-1)) &= a_{10} + a_{11}T(n-1) = a_{10} + a_{11} \frac{T(n) - \tanh(c_1)}{1 - T(n) \tanh(c_1)} \\ U_{n+1}(T(n+1)) &= a_{10} + a_{11}T(n+1) = a_{10} + a_{11} \frac{T(n) + \tanh(c_1)}{1 + T(n) \tanh(c_1)}. \end{aligned}$$

## Step 3:

- Determine the algebraic system for the unknown coefficients  $a_{ij}$  by balancing the coefficients of the various powers of  $T(n)$ .
- Example: Algebraic system for Toda lattice

$$c_2^2 - \tanh^2(c_1) - a_{11}c_2 \tanh^2(c_1) = 0, \quad c_2 - a_{11} = 0$$

#### Step 4:

- Solve the nonlinear algebraic system with parameters.  
Reject complex solutions? Test the solutions.
- Example: Solution of algebraic system for Toda lattice

$$a_{10} = \text{free}, \quad a_{11} = c_2 = \pm \sinh(c_1)$$

#### Step 5:

- Return to the original variables.  
Test the final solution in the original equations
- Example: Solitary wave solution for Toda lattice:

$$u_n(t) = a_0 \pm \sinh(c_1) \tanh [c_1 n \pm \sinh(c_1) t + \delta].$$

## Example: System of DDEs: Relativistic Toda lattice

$$\begin{aligned}\dot{u}_n &= (1 + \alpha u_n)(v_n - v_{n-1}) \\ \dot{v}_n &= v_n(u_{n+1} - u_n + \alpha v_{n+1} - \alpha v_{n-1}).\end{aligned}$$

**Step 1:** Change of variables

$$u_n(x, t) = U_n(T(n)), \quad v_n(x, t) = V_n(T(n)),$$

with

$$T(n) = \tanh [c_1 n + c_2 t + \delta] = \tanh \xi.$$

gives

$$\begin{aligned}c_2(1 - T^2)U'_n - (1 + \alpha U_n)(V_n - V_{n-1}) &= 0 \\ c_2(1 - T^2)V'_n - V_n(U_{n+1} - U_n + \alpha V_{n+1} - \alpha V_{n-1}) &= 0.\end{aligned}$$

**Step 2:** Seek polynomial solutions

$$\begin{aligned}U_n(T(n)) &= \sum_{j=0}^{M_1} a_{1j} T(n)^j \\ V_n(T(n)) &= \sum_{j=0}^{M_2} a_{2j} T(n)^j.\end{aligned}$$

Balance the highest power terms in  $T(n)$  to determine  $M_1$ , and  $M_2$  :

$$M_1 + 1 = M_1 + M_2, \quad M_2 + 1 = M_1 + M_2$$

gives  $M_1 = M_2 = 1$ .

Hence,

$$U_n = a_{10} + a_{11}T(n), \quad V_n = a_{20} + a_{21}T(n).$$

**Step 3:** Algebraic system for  $a_{ij}$  :

$$\begin{aligned}
 -a_{11} c_2 + a_{21} \tanh(c_1) + \alpha a_{10} a_{21} \tanh(c_1) &= 0 \\
 a_{11} \tanh(c_1) (\alpha a_{21} + c_2) &= 0 \\
 -a_{21} c_2 + a_{11} a_{20} \tanh(c_1) + 2\alpha a_{20} a_{21} \tanh(c_1) &= 0 \\
 \tanh(c_1) (a_{11} a_{21} + 2\alpha a_{21}^2 - a_{11} a_{20} \tanh(c_1)) &= 0 \\
 a_{21} \tanh^2(c_1) (c_2 - a_{11}) &= 0
 \end{aligned}$$

**Step 4:** Solution of the algebraic system

$$a_{10} = -c_2 \coth(c_1) - \frac{1}{\alpha}, \quad a_{11} = c_2, \quad a_{20} = \frac{c_2 \coth(c_1)}{\alpha}, \quad a_{21} = -\frac{c_2}{\alpha}.$$

**Step 5:** Solitary wave solution in original variables:

$$\begin{aligned}
 u_n(t) &= -c_2 \coth(c_1) - \frac{1}{\alpha} + c_2 \tanh [c_1 n + c_2 t + \delta] \\
 v_n(t) &= \frac{c_2 \coth(c_1)}{\alpha} - \frac{c_2}{\alpha} \tanh [c_1 n + c_2 t + \delta].
 \end{aligned}$$

Solving/Analyzing Systems of Algebraic Equations with Parameters  
 Class of fifth-order evolution equations with parameters:

$$u_t + \alpha\gamma^2 u^2 u_x + \beta\gamma u_x u_{2x} + \gamma u u_{3x} + u_{5x} = 0.$$

### Well-Known Special cases

Lax case:  $\alpha = \frac{3}{10}, \beta = 2, \gamma = 10$ . **Two** solutions:

$$u(x, t) = 4c_1^2 - 6c_1^2 \tanh^2 [c_1 x - 56c_1^5 t + \delta]$$

and

$$u(x, t) = a_0 - 2c_1^2 \tanh^2 [c_1 x - 2(15a_0^2 c_1 - 40a_0 c_1^3 + 28c_1^5)t + \delta]$$

where  $a_0$  is arbitrary.

Sawada-Kotera case:  $\alpha = \frac{1}{5}, \beta = 1, \gamma = 5$ . **Two** solutions:

$$u(x, t) = 8c_1^2 - 12c_1^2 \tanh^2 [c_1 x - 16c_1^5 t + \delta]$$

and

$$u(x, t) = a_0 - 6c_1^2 \tanh^2 [c_1 x - (5a_0^2 c_1 - 40a_0 c_1^3 + 76c_1^5)t + \delta]$$

where  $a_0$  is arbitrary.

Kaup-Kupershmidt case:  $\alpha = \frac{1}{5}, \beta = \frac{5}{2}, \gamma = 10$ . **Two** solutions:

$$u(x, t) = c_1^2 - \frac{3}{2}c_1^2 \tanh^2 [c_1 x - c_1^5 t + \delta]$$

and

$$u(x, t) = 8c_1^2 - 12c_1^2 \tanh^2 [c_1 x - 176c_1^5 t + \delta],$$

no free constants!

Ito case:  $\alpha = \frac{2}{9}, \beta = 2, \gamma = 3$ . **One** solution:

$$u(x, t) = 20c_1^2 - 30c_1^2 \tanh^2 [c_1 x - 96c_1^5 t + \delta].$$

## What about the General case?

Q1: Can we retrieve the special solutions?

Q2: What are the condition(s) on the parameters  $\alpha, \beta, \gamma$  for solutions of tanh-type to **exist**?

Tanh solutions:

$$u(x, t) = a_0 + a_1 \tanh [c_1 x + c_2 t + \delta] + a_2 \tanh^2 [c_1 x + c_2 t + \delta].$$

Nonlinear algebraic system must be analyzed, solved (or reduced!):

$$a_1(\alpha\gamma^2 a_2^2 + 6\gamma a_2 c_1^2 + 2\beta\gamma a_2 c_1^2 + 24c_1^4) = 0$$

$$a_1(\alpha\gamma^2 a_1^2 + 6\alpha\gamma^2 a_0 a_2 + 6\gamma a_0 c_1^2 - 18\gamma a_2 c_1^2 - 12\beta\gamma a_2 c_1^2 - 120c_1^4) = 0$$

$$\alpha\gamma^2 a_2^2 + 12\gamma a_2 c_1^2 + 6\beta\gamma a_2 c_1^2 + 360c_1^4 = 0$$

$$2\alpha\gamma^2 a_1^2 a_2 + 2\alpha\gamma^2 a_0 a_2^2 + 3\gamma a_1^2 c_1^2 + \beta\gamma a_1^2 c_1^2 + 12\gamma a_0 a_2 c_1^2 - 8\gamma a_2^2 c_1^2 - 8\beta\gamma a_2^2 c_1^2 - 480a_2 c_1^4 = 0$$

$$a_1(\alpha\gamma^2 a_0^2 c_1 - 2\gamma a_0 c_1^3 + 2\beta\gamma a_2 c_1^3 + 16c_1^5 + c_2) = 0$$

$$\alpha\gamma^2 a_0 a_1^2 c_1 + \alpha\gamma^2 a_0^2 a_2 c_1 - \gamma a_1^2 c_1^3 - \beta\gamma a_1^2 c_1^3 - 8\gamma a_0 a_2 c_1^3 + 2\beta\gamma a_2^2 c_1^3 + 136a_2 c_1^5 + a_2 c_2 = 0$$

Unknowns:  $a_0, a_1, a_2$ .

Parameters:  $c_1, c_2, \alpha, \beta, \gamma$ .

**Solve** and **Reduce** cannot be used on the whole system!

## Strategy to Solve/Reduce Nonlinear Systems

Assumptions:

- All  $c_i \neq 0$
- Parameters ( $\alpha, \beta, \gamma, \dots$ ) are nonzero. Otherwise the highest powers  $M_i$  may change.
- All  $a_{j M_i} \neq 0$ . Coefficients of highest power in  $U_i$  are present.
- Solve for  $a_{ij}$ , then  $c_i$ , then find conditions on parameters.

Strategy followed by hand:

- Solve all linear equations in  $a_{ij}$  first (cost: branching). Start with the ones without parameters. Capture constraints in the process.
- Solve linear equations in  $c_i$  if they are free of  $a_{ij}$ .
- Solve linear equations in parameters if they free of  $a_{ij}, c_i$ .
- Solve quasi-linear equations for  $a_{ij}, c_i$ , parameters.
- Solve quadratic equations for  $a_{ij}, c_i$ , parameters.
- Eliminate cubic terms for  $a_{ij}, c_i$ , parameters, without solving.
- Show remaining equations, if any.

Alternatives:

- Use (adapted) Gröbner Basis Techniques.
- Use combinatorics on coefficients  $a_{ij} = 0$  or  $a_{ij} \neq 0$ .

**Actual Solution:** Two major cases:

CASE 1:  $a_1 = 0$ , two subcases

**Subcase 1-a:**

$$a_2 = -\frac{3}{2}a_0$$

$$c_2 = c_1^3(24c_1^2 - \beta\gamma a_0)$$

where  $a_0$  is one of the two roots of the quadratic equation:

$$\alpha\gamma^2 a_0^2 - 8\gamma a_0 c_1^2 - 4\beta\gamma a_0 c_1^2 + 160c_1^4 = 0.$$

**Subcase 1-b:** If  $\beta = 10\alpha - 1$ , then

$$a_2 = -\frac{6}{\alpha\gamma}c_1^2$$

$$c_2 = -\frac{1}{\alpha}(\alpha^2\gamma^2 a_0^2 c_1 - 8\alpha\gamma a_0 c_1^3 + 12c_1^5 + 16\alpha c_1^5)$$

where  $a_0$  is arbitrary.

CASE 2:  $a_1 \neq 0$ , then

$$\alpha = \frac{1}{392}(39 + 38\beta + 8\beta^2)$$

and

$$a_2 = -\frac{168}{\gamma(3 + 2\beta)}c_1^2$$

provided  $\beta$  is one of the roots of

$$(104\beta^2 + 886\beta + 1487)(520\beta^3 + 2158\beta^2 - 1103\beta - 8871) = 0$$

**Subcase 2-a:** If  $\beta^2 = -\frac{1}{104}(886\beta + 1487)$ , then

$$\alpha = -\frac{2\beta + 5}{26}$$

$$a_0 = -\frac{49c_1^2(9983 + 4378\beta)}{26\gamma(8 + 3\beta)(3 + 2\beta)^2}$$

$$a_1 = \pm\frac{336c_1^2}{\gamma(3 + 2\beta)}$$

$$a_2 = -\frac{168c_1^2}{\gamma(3 + 2\beta)}$$

$$c_2 = -\frac{364c_1^5(3851 + 1634\beta)}{6715 + 2946\beta}.$$

**Subcase 2-b:** If  $\beta^3 = \frac{1}{520}(8871 + 1103\beta - 2158\beta^2)$ , then

$$\alpha = \frac{39 + 38\beta + 8\beta^2}{392}$$

$$a_0 = \frac{28c_1^2(6483 + 5529\beta + 1066\beta^2)}{(3 + 2\beta)(23 + 6\beta)(81 + 26\beta)\gamma}$$

$$a_1^2 = \frac{28224c_1^4(4\beta - 1)(26\beta - 17)}{(3 + 2\beta)^2(23 + 6\beta)(81 + 26\beta)\gamma^2}$$

$$a_2 = -\frac{168c_1^2}{\gamma(3 + 2\beta)}$$

$$c_2 = -\frac{8c_1^5(1792261977 + 1161063881\beta + 188900114\beta^2)}{959833473 + 632954969\beta + 105176786\beta^2}.$$

## Implementation Issues – Software Demo – Future Work

- Demonstration of Mathematica package for tanh/sech methods.
- Long term goal: Develop PDESolve for closed form solutions of nonlinear PDEs and DDEs.
- Implement various methods: Lie symmetry (similarity) methods.
- Look at other types of explicit solutions involving
  - hyperbolic functions sinh, cosh, tanh, ...
  - other special functions.
  - complex exponentials combined with sech or tanh.

- Example: Set of ODEs from quantum field theory

$$\begin{aligned} u_{xx} &= -u + u^3 + av^2 \\ v_{xx} &= bv + cv^3 + av(u^2 - 1). \end{aligned}$$

Try solutions ( $c_2 = 0$  for ODEs)

$$u_i(x, t) = \sum_{j=0}^{M_i} a_{ij} \tanh^j[c_1x + c_2t + \delta] + \sum_{j=0}^{N_i} b_{ij} \operatorname{sech}^{2j+1}[c_1x + c_2t + \delta].$$

or

$$u_i(x, t) = \sum_{j=0}^{M_i} (\tilde{a}_{ij} + \tilde{b}_{ij} \operatorname{sech}[c_1x + c_2t + \delta]) \tanh^j[c_1x + c_2t + \delta].$$

Solitary wave solutions:

$$\begin{aligned} u &= \pm \tanh\left[\sqrt{\frac{a^2 - c}{2(a - c)}}x + \delta\right] \\ v &= \pm \sqrt{\frac{1 - a}{a - c}} \operatorname{sech}\left[\sqrt{\frac{a^2 - c}{2(a - c)}}x + \delta\right], \end{aligned}$$

provided  $b = \sqrt{\frac{a^2 - c}{2(a - c)}}$ .

- Example: Nonlinear Schrödinger equation (focusing/defocusing):

$$i u_t + u_{xx} \pm |u|^2 u = 0.$$

Bright soliton solution (+ sign):

$$u(x, t) = \frac{k}{\sqrt{2}} \exp\left[i\left(\frac{c}{2}x + \left(k^2 - \frac{c^2}{4}\right)t\right)\right] \operatorname{sech}[k(x - ct - x_0)]$$

Dark soliton solution (− sign):

$$u(x, t) = \frac{1}{\sqrt{2}} \exp\left[i\left(Kx - \left(2k^2 + 3K^2 - 2Kc + \frac{c^2}{2}\right)t\right)\right] \left\{ k \tanh[k(x - ct - x_0)] - i\left(K - \frac{c}{2}\right) \right\}.$$

- Example: Nonlinear sine-Gordon equation (light cone coordinates):

$$u_{xt} = \sin u.$$

Setting  $\Phi = u_x$ ,  $\Psi = \cos(u) - 1$ , gives

$$\begin{aligned} \Phi_{xt} - \Phi - \Phi\Psi &= 0 \\ 2\Psi + \Psi^2 + \Phi_t^2 &= 0. \end{aligned}$$

Solitary wave solution (kink):

$$\begin{aligned} \Phi &= u_x = \pm \frac{1}{\sqrt{-c}} \operatorname{sech}\left[\frac{1}{\sqrt{-c}}(x - ct) + \delta\right], \\ \Psi &= \cos(u) - 1 = 1 - 2 \operatorname{sech}^2\left[\frac{1}{\sqrt{-c}}(x - ct) + \delta\right], \end{aligned}$$

in final form:

$$u(x, t) = \pm 4 \arctan\left(\exp\left(\frac{1}{\sqrt{-c}}(x - ct) + \delta\right)\right).$$

- Example: Coupled nonlinear Schrödinger equations:

$$\begin{aligned}i u_t &= u_{xx} + u(|u|^2 + h|v|^2) \\i v_t &= v_{xx} + v(|v|^2 + h|u|^2)\end{aligned}$$

Seek particular solutions

$$\begin{aligned}u(x, t) &= a \tanh(\mu x) \exp(iAt) \\v(x, t) &= b \operatorname{sech}(\mu x) \exp(iBt).\end{aligned}$$

- Seek solutions  $u(x, t) = U(F(\xi))$ , where derivatives of  $F(\xi)$  are polynomial in  $F$ .

Now,

$$F'(\xi) = 1 - F^2(\xi) \quad \longrightarrow \quad F = \tanh(\xi).$$

Other choices are possible.

- Add the constraining differential equations to the system of PDEs directly.
- Why are  $\tanh$  and  $\operatorname{sech}$  solutions so prevalent?
- Other applications:

Computation of conservation laws, symmetries, first integrals, etc. leading to **linear** parameterized systems for unknowns coefficients (see *InvariantsSymmetries* by Göktaş and Hereman).